



Intra- and Inter-molecular Hydrogen Bonding in Overtone Spectroscopy

IUPAC workshop, Pisa, September 5-9, 2005

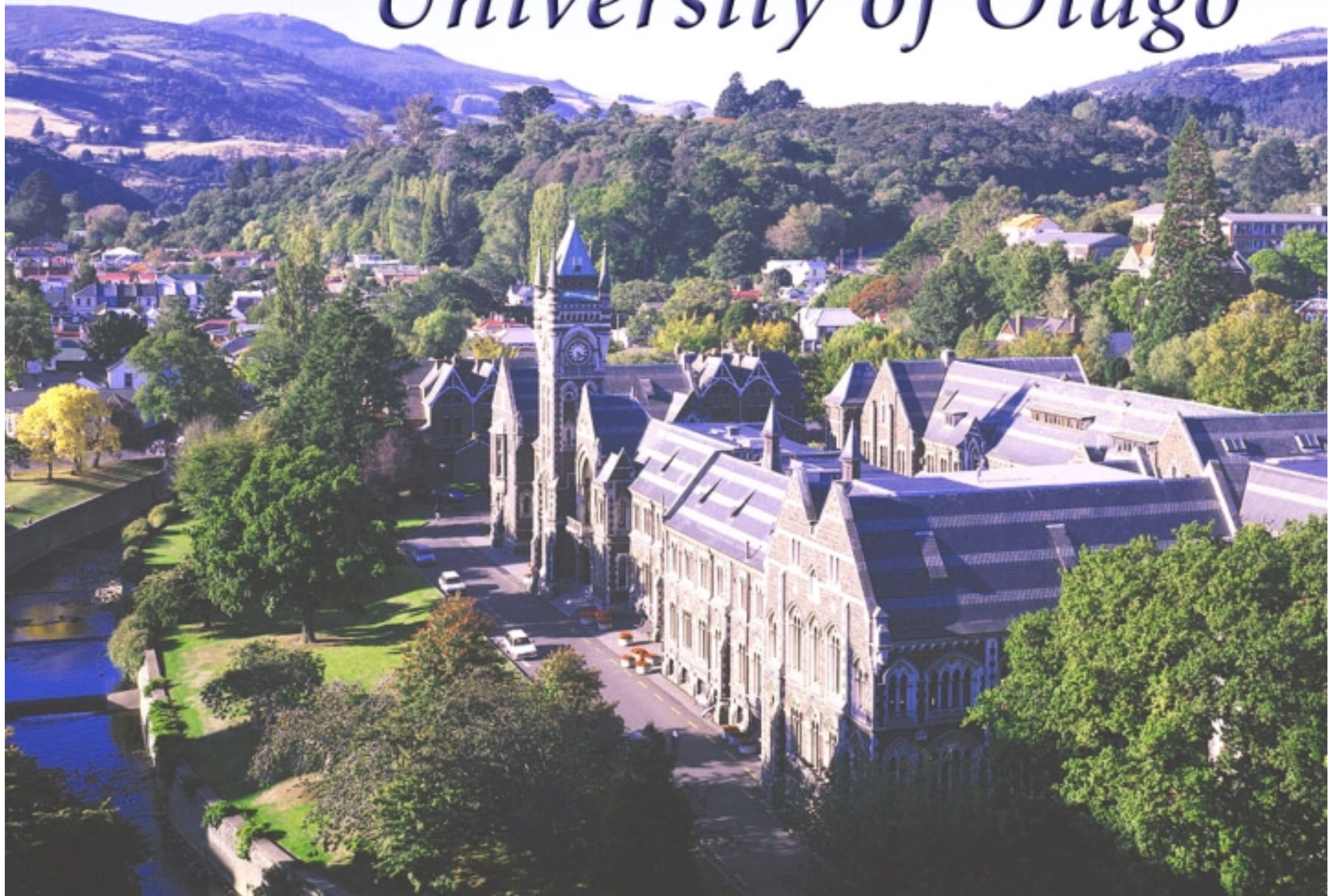


Henrik G. Kjaergaard

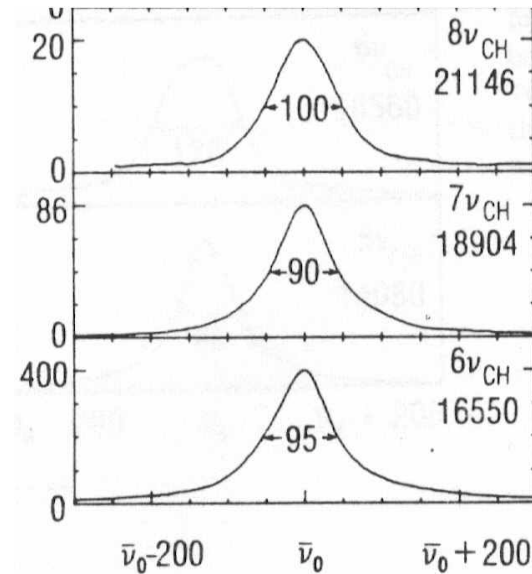
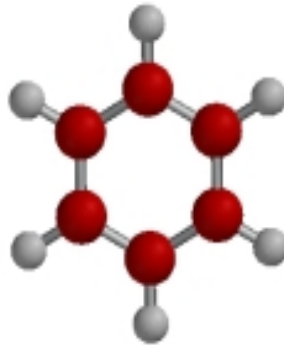
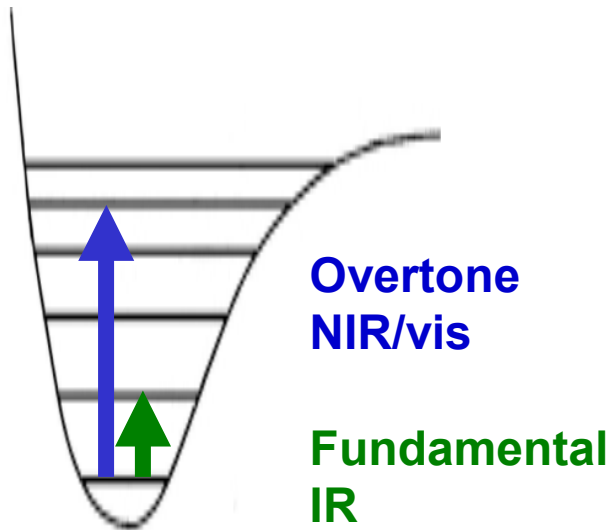
Department of Chemistry,
University of Otago,
Dunedin, New Zealand.



University of Otago



Vibrational Overtones



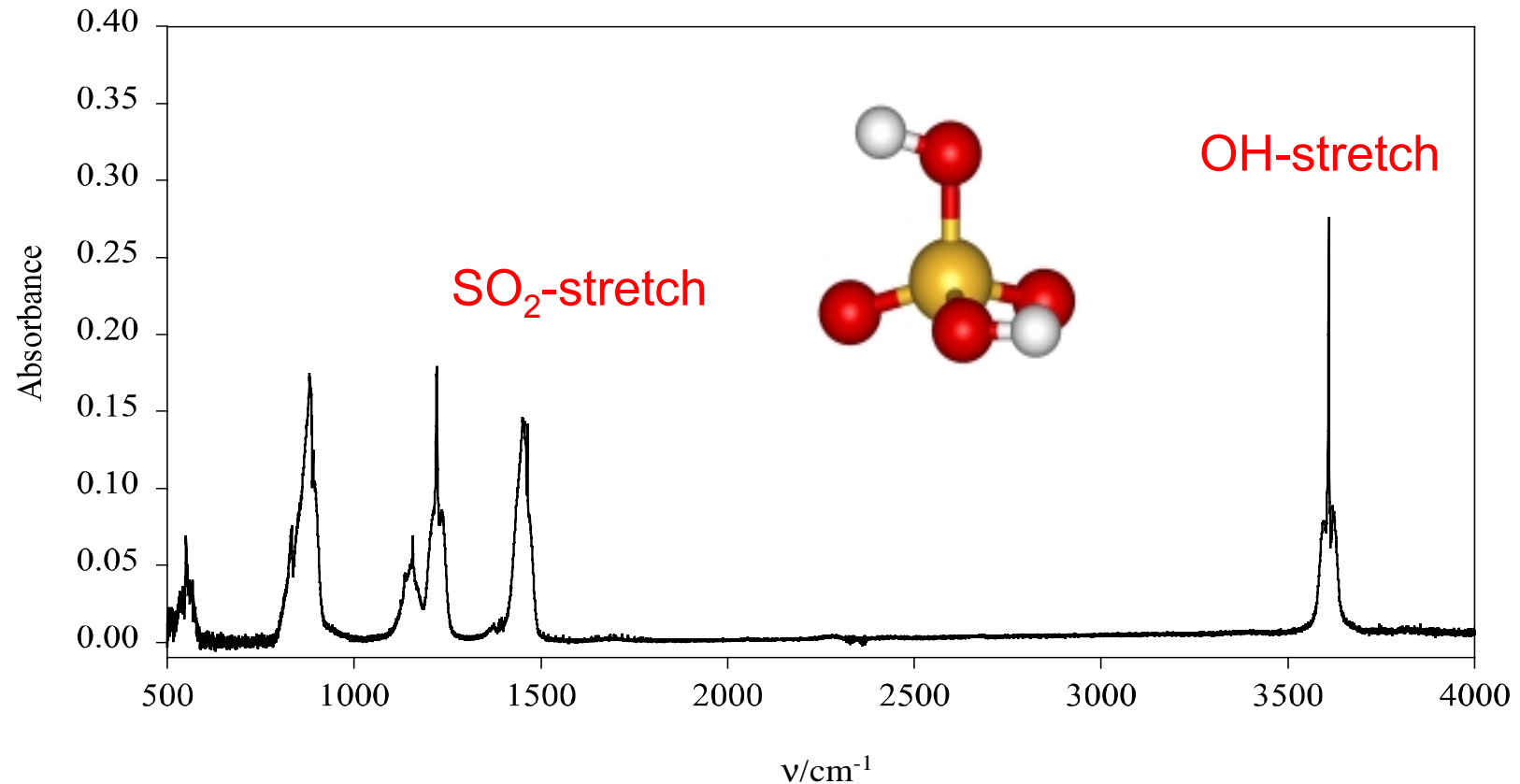
Overtone spectra are dominated by XH-stretching vibrations.

Described well by Local modes.

Vibrational energy is localized in a single bond or in a linear combination of equivalent bonds.

Reddy *et al*, J. Chem.Phys.
76, 2814 (1982).

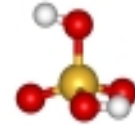
IR Spectrum of Sulfuric Acid Vapor



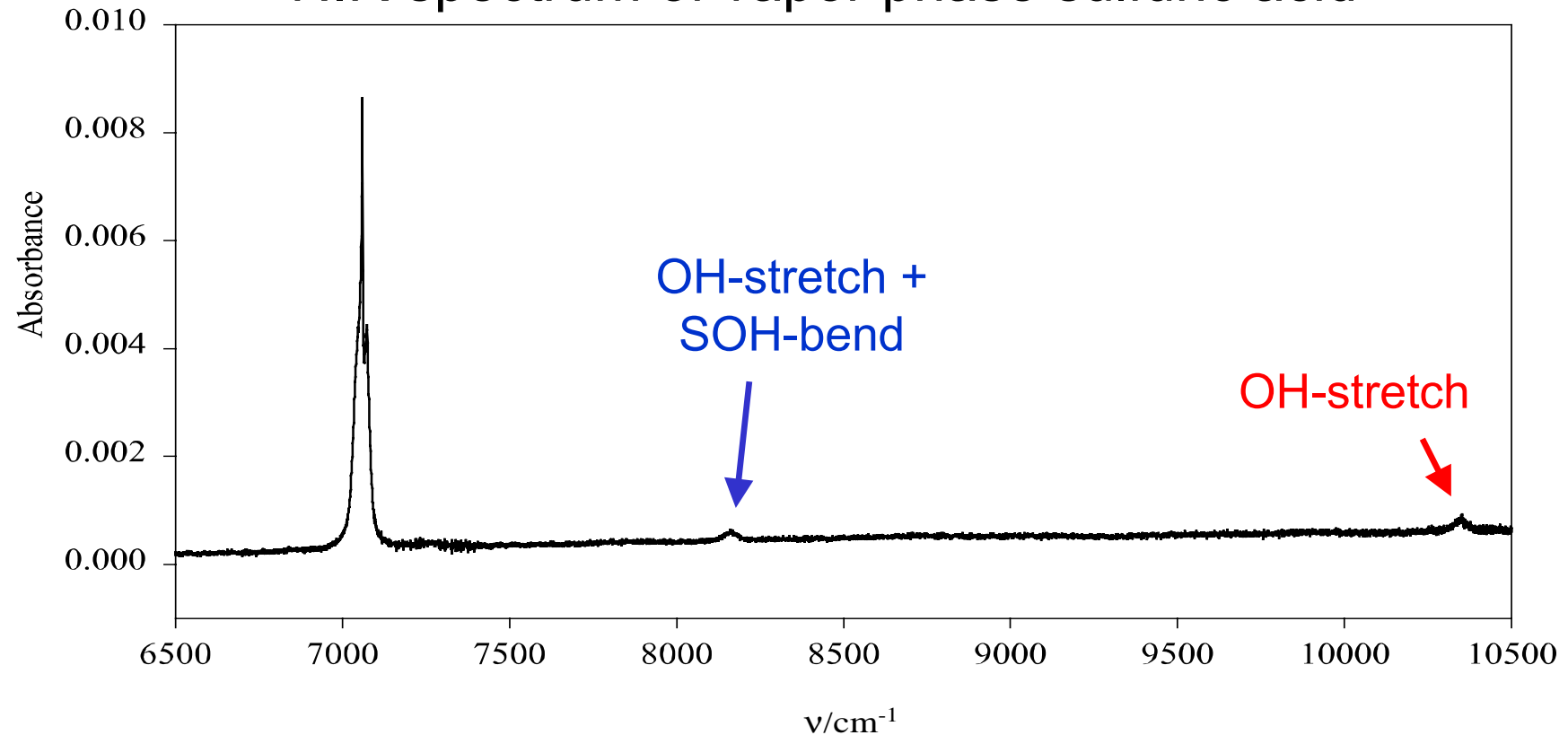
Several normal modes of vibration seen in the fundamental region.

Hintze, Kjaergaard, Vaida, Burkholder, JPC A 2003.

Normal versus Local modes



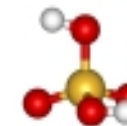
NIR spectrum of vapor phase sulfuric acid



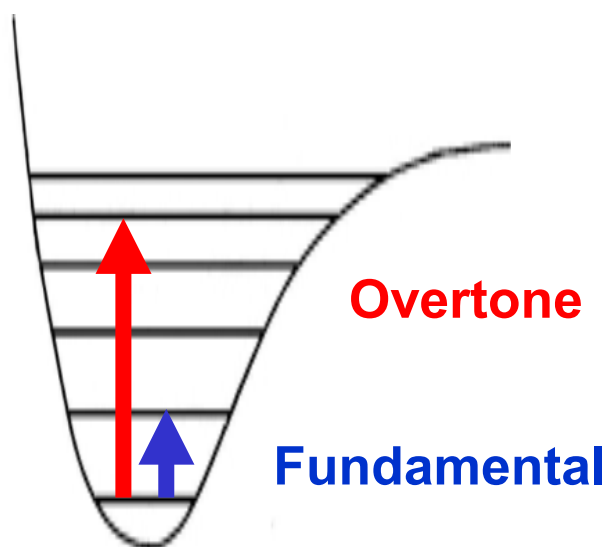
Overtone spectra dominated by OH-stretching vibrations

Described well by anharmonic local modes.

OH-stretching frequencies



Vibrational energies described well by isolated OH bond.



Δv_{OH}	Freq. (cm^{-1})	
	Obs.	Calc.
1	3609.2	3609.2
2	7060.7	7059.4
3	10350.3	10350.6
4		13483

Best fit Morse parameters $\tilde{\omega} = 3768.2 \text{ cm}^{-1}$ $\tilde{\omega}_x = 79.5 \text{ cm}^{-1}$

Vibrational local modes

OH-stretch as anharmonic (Morse) oscillator:

$$(H - E_{0v})/hc = v\tilde{\omega} - (v^2 + v)\tilde{\omega}x \Rightarrow E_{vib} \text{ and } \psi_{vib}$$

$\tilde{\omega}$ and $\tilde{\omega}x$ from observed peak positions

Dimensionless oscillator strength:

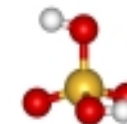
$$f_{eg} \propto \tilde{\nu}_{eg} \left| \langle e | \vec{\mu} | g \rangle \right|^2$$

Dipole moment function of single OH bond

$$\vec{\mu}(q) = \sum_i \vec{\mu}_i q^i$$

From *ab initio* calculated dipole grid

OH-stretching intensities



Use 1D Morse oscillator model with:

Δv_{OH}	Int. (km/mol)	
	Obs.	Calc.
1	180	205
2	6.3	7.0
3	0.23	0.29
4		0.014

$$\tilde{\omega} = 3768.2 \text{ cm}^{-1}$$

$$\tilde{\omega}_x = 79.5 \text{ cm}^{-1}$$

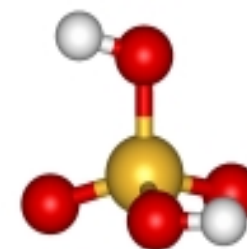
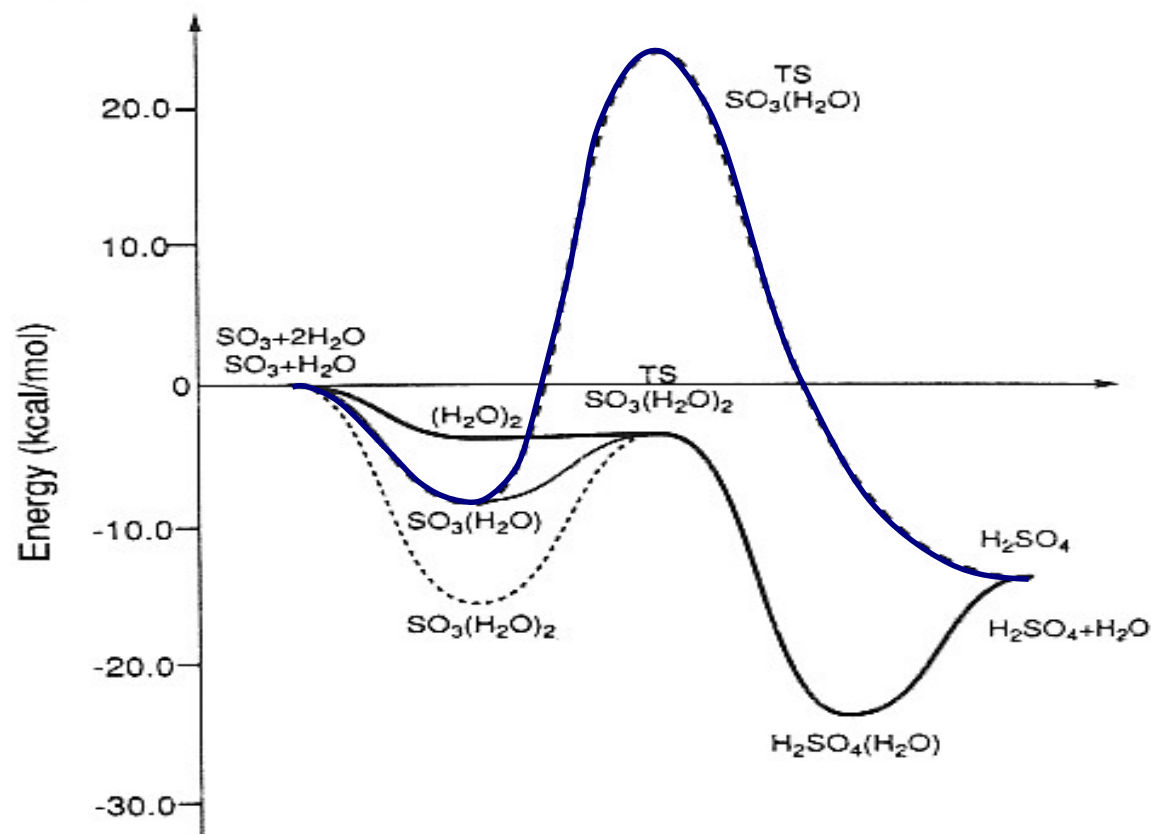
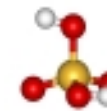
B3LYP/6-311++G(2d,2p)
Dipole Moment Function

Adding SOH-bending mode
has little effect on OH-stretch

G98 FREQ Calculation with B3LYP/6-311++G(2d,2p)

Harmonic Oscillator Linear Dipole calculation \Rightarrow 252 km/mol

Photolysis of H₂SO₄



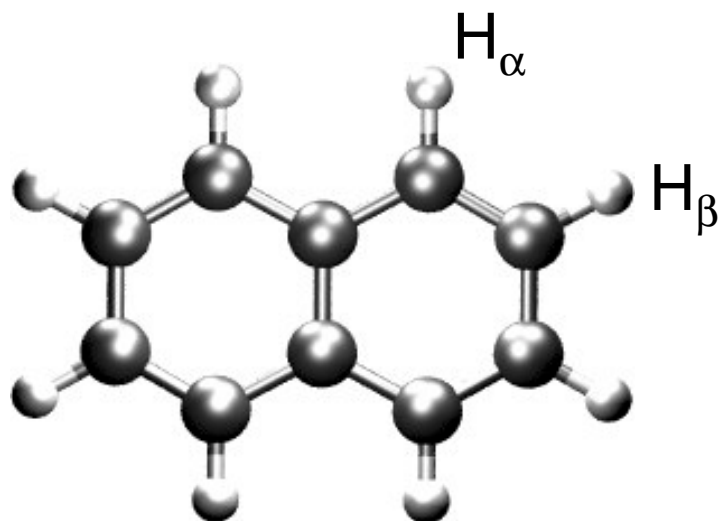
Thus H₂SO₄ photolysis can proceed via overtone excitations.

Morokuma, Mugurama, JACS (1994).

Vaida, Kjaergaard, Hintze, Donaldson, Science (2003).

Relative intensities

Naphthalene

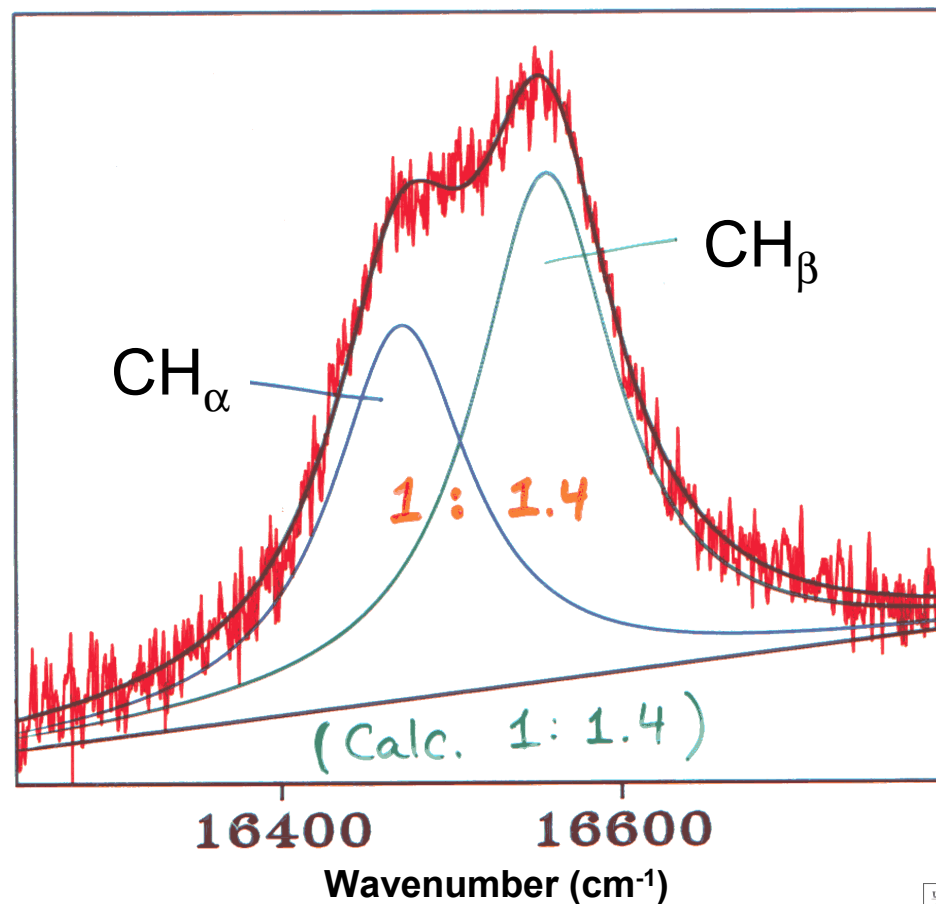


HF/6-311+G(d,p)
calculated bond lengths:

$$CH_{\alpha} = 1.085 \text{ \AA}$$

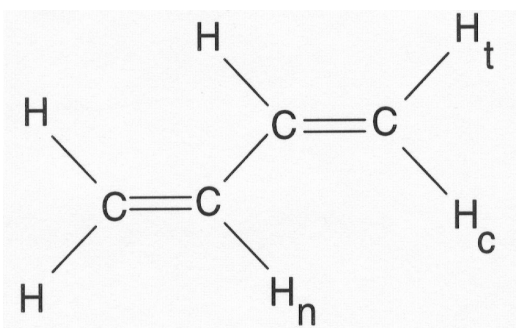
$$CH_{\beta} = 1.084 \text{ \AA}$$

Vapor phase $\Delta\nu_{CH} = 6$ spectrum



Relative intensities

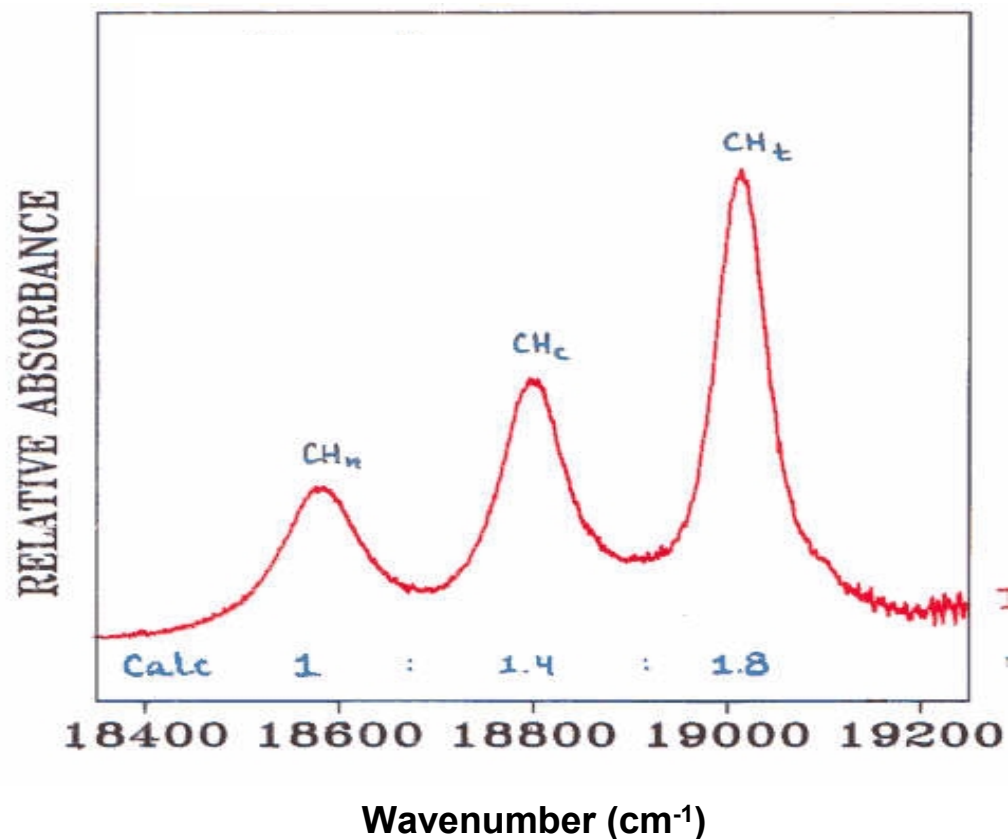
1,3-butadiene



Calculate *ab initio* DMF for each of the three non-equivalent CH bonds.

HF/6-31G(d) method.

Vapor phase $\Delta\nu_{\text{CH}} = 7$ spectrum

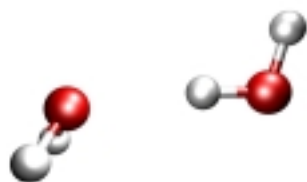


Kjaergaard, Turnbull, Henry, JPC 1993.

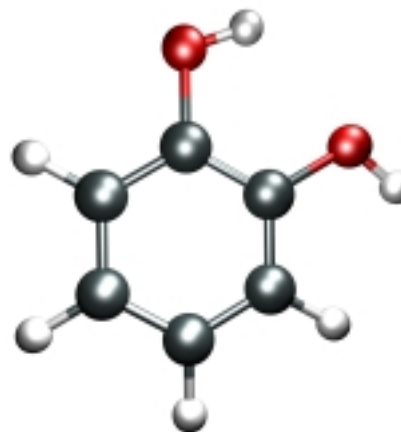
Inter- vs Intra-molecular interaction

How are overtone spectra effected by small change from weak molecular interactions?

Weak van der Waals interaction (including hydrogen bonding) can occur both inter-molecular as in complexes and intra-molecular.



Water dimer

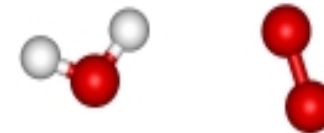
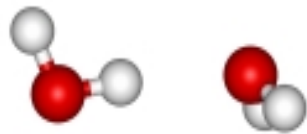


Catechol

Complexes in the Atmosphere

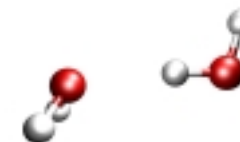
$O_2 \cdot O_2$ and $O_2 \cdot N_2$ complexes have been shown to absorb about 1 W/m^2 of incoming solar radiation.

Hydrated complexes, $H_2O \cdot X$, are likely to contribute.

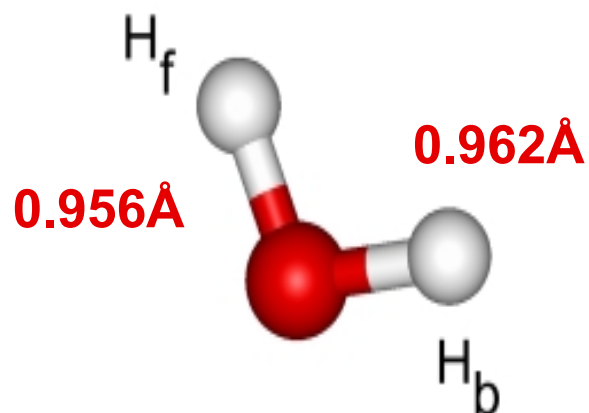


Contribution depends on position, intensity and shape of spectroscopic transitions as well as atmospheric abundance.

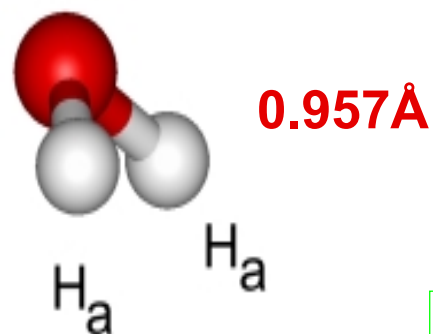
Water dimer, $\text{H}_2\text{O}\cdot\text{H}_2\text{O}$



QCISD/6-311++G(2d,2p) optimized geometry



H donor

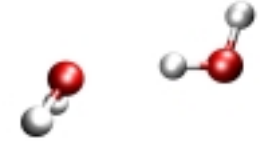


H acceptor

Water monomer,
 $R_{\text{OH}} = 0.958 \text{ \AA}$

OH bond involved in hydrogen bonding is significantly longer \Rightarrow frequency red shift

Water dimer, H₂O•H₂O



Simple vibrational model for water dimer

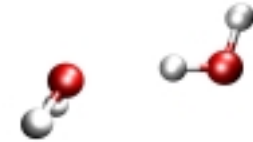
$$\hat{H} = \hat{H}_{acceptor} + \hat{H}_{donor}$$

Each H₂O unit is modeled by two OH-stretching and one HOH-bending local mode oscillator.

We use the Harmonically Coupled Anharmonic Oscillators (HCAO) local mode model for each of the H₂O units.

Not sufficient experimental data to determine local mode parameters. Obtain these from ab initio calculations.

Local mode parameters



Local mode parameters for water

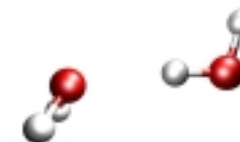
	QCISD	Expt.	Scaling factor
$\tilde{\omega}$	3934	3870	0.9836
$\tilde{\omega}_x$	96.5	82.1	0.851

QCISD calculated Local mode parameters for water dimer after scaling

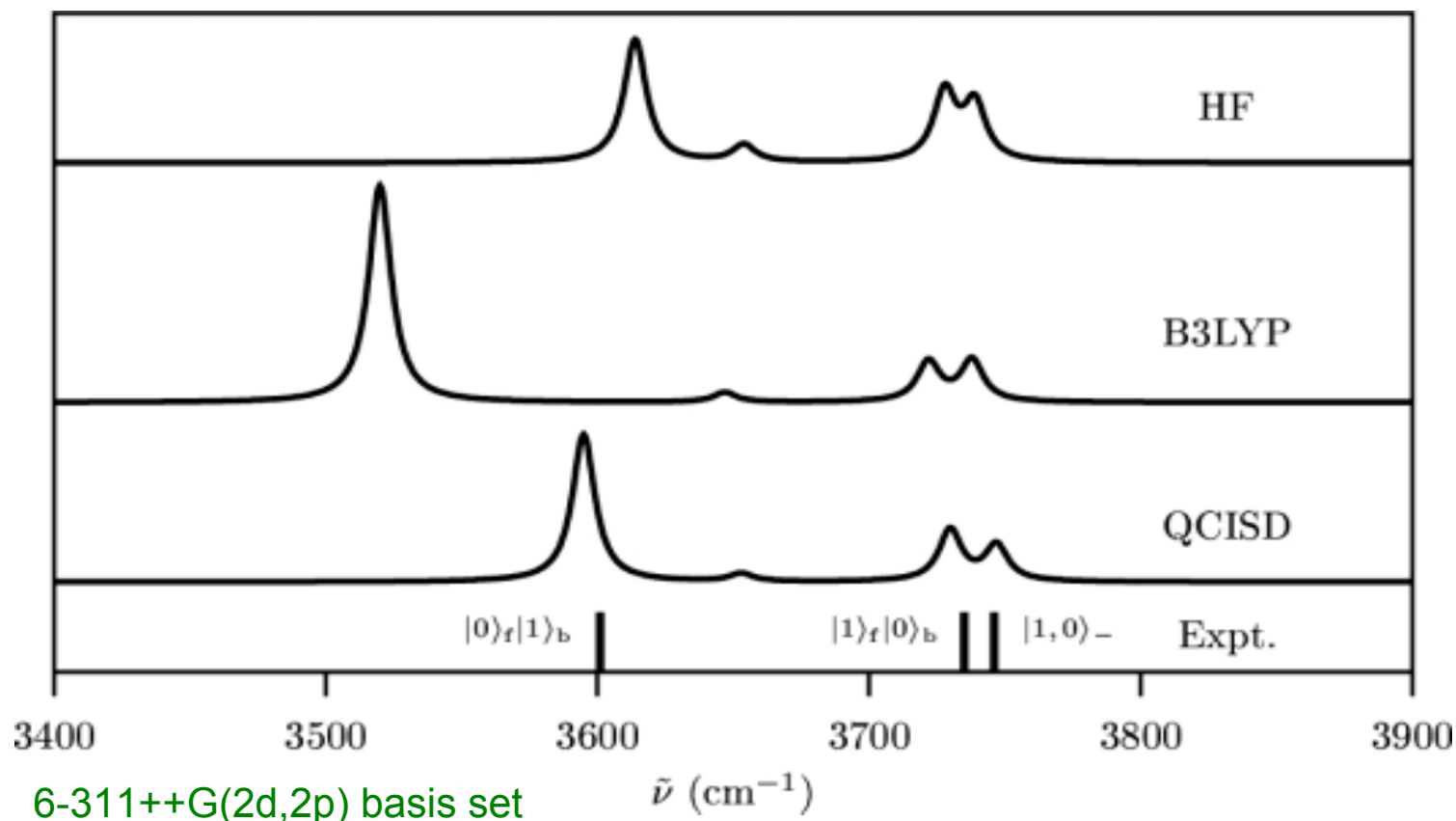
	OH _a	OH _f	OH _b	H ₂ O expt.
$\tilde{\omega}$	3863	3878	3781	3870
$\tilde{\omega}_x$	81.7	82.3	85.4	82.1

Low, Kjaergaard, JCP 1999.

Ab initio methods



OH-stretching fundamental transitions in $\text{H}_2\text{O}\cdot\text{H}_2\text{O}$



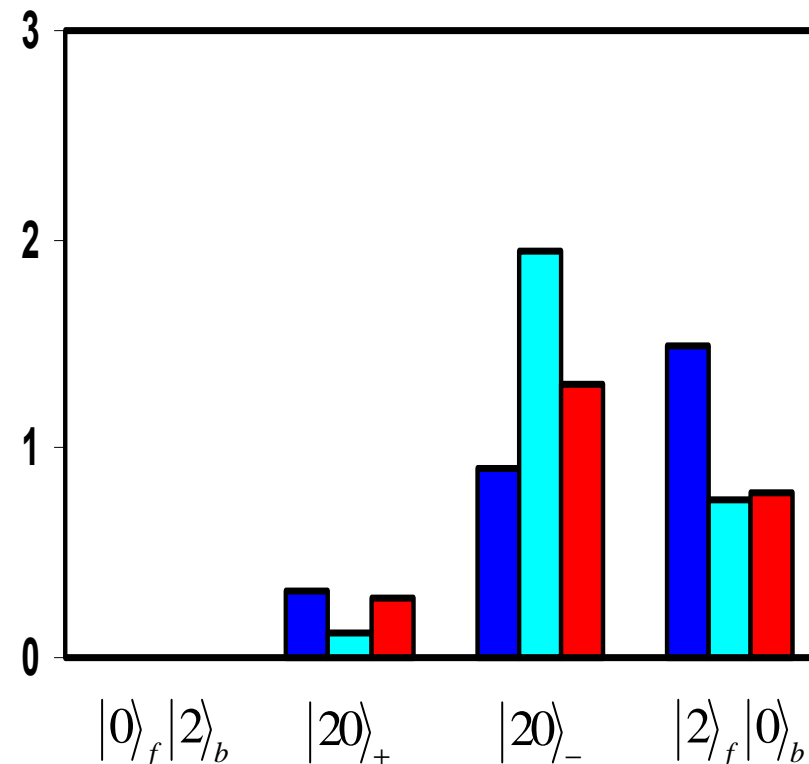
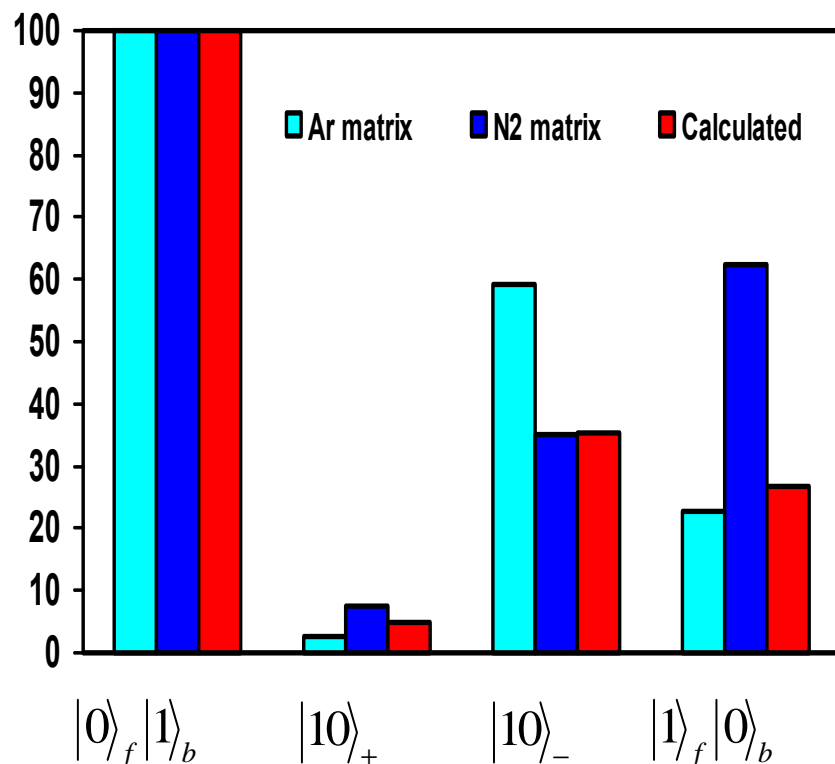
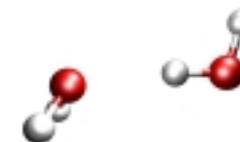
6-311++G(2d,2p) basis set

Low, Kjaergaard, JCP 1999.

Huang, Miller, JCP 1989.

Huisken, Kaloudis, Kulcke, JCP 1996.

Water dimer, OH-stretching



First overtone of OH_b -stretching transition is very weak, only recently observed in Ne matrix.

Perchard, Chem. Phys. (2001)

Schofield, Kjaergaard, PCCP 2003.

Cancellation of terms

Expansion: $\langle v | \mu | 0 \rangle \approx \mu_1 \langle v | q | 0 \rangle + \mu_2 \langle v | q^2 | 0 \rangle + \dots$

Matrix elements for a typical OH-stretching Morse oscillator

Fundamental: $\langle 1 | q | 0 \rangle \approx 0.07$ $\langle 1 | q^2 | 0 \rangle \approx 0.0035$

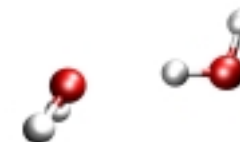
1st overtone: $\langle 2 | q | 0 \rangle \approx -0.0072$ $\langle 2 | q^2 | 0 \rangle \approx 0.0062$

Ab initio dipole expansion coefficients

OH bond in H₂O: $\mu_1 \approx 0.76$ $\mu_2 \approx -0.60$

OH_b bond in (H₂O)₂: $\mu_1 \approx 2.53$ $\mu_2 \approx 2.66$

Water dimer, OH-stretching



OH-stretching frequencies in the $\Delta v_{\text{OH}} = 2$ region (in cm^{-1})

Mode	Ne Matrix	Jet-cooled	HCAO calc.
$ 0\rangle_f 2\rangle_b$	7018	-	7044
$ 20\rangle_+$	7207	7193	7198
$ 2\rangle_f 0\rangle_b$	7237	7240	7236
$ 20\rangle_-$	7245	7250	7238
$ 1\rangle_f 1\rangle_b$	7362	-	7370

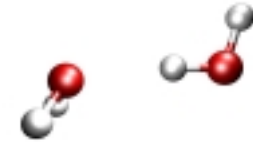
Bouteiller, Perchard, Chem. Phys. (2004)

Nizkorodov, Ziemkiewicz, Nesbitt, Knight,
JCP (2005)

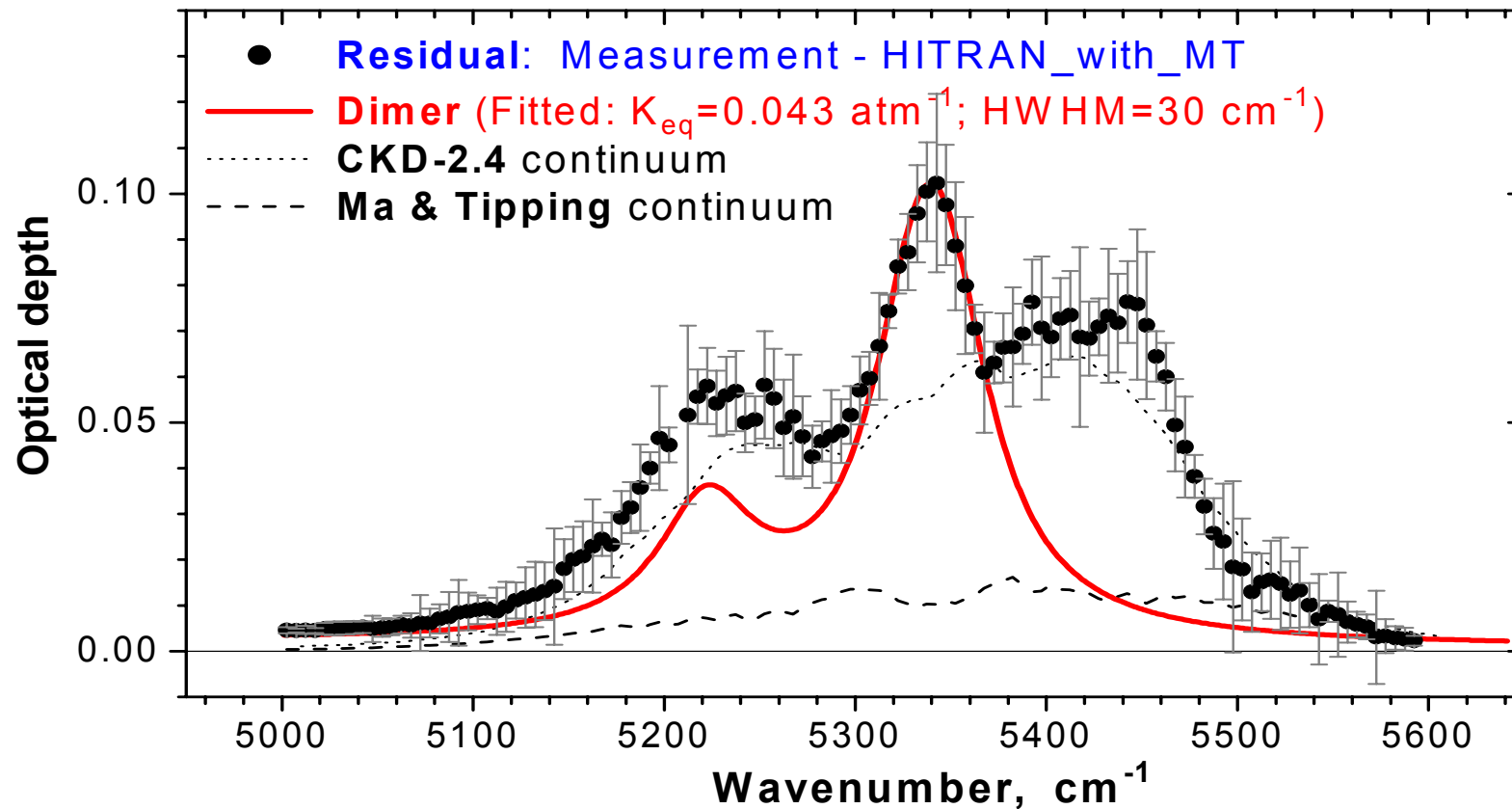
Schofield, Kjaergaard, PCCP 2003.

HCAO calculation at
QCISD/6-311++G(2d,2p)

Water dimer, stretch-bend



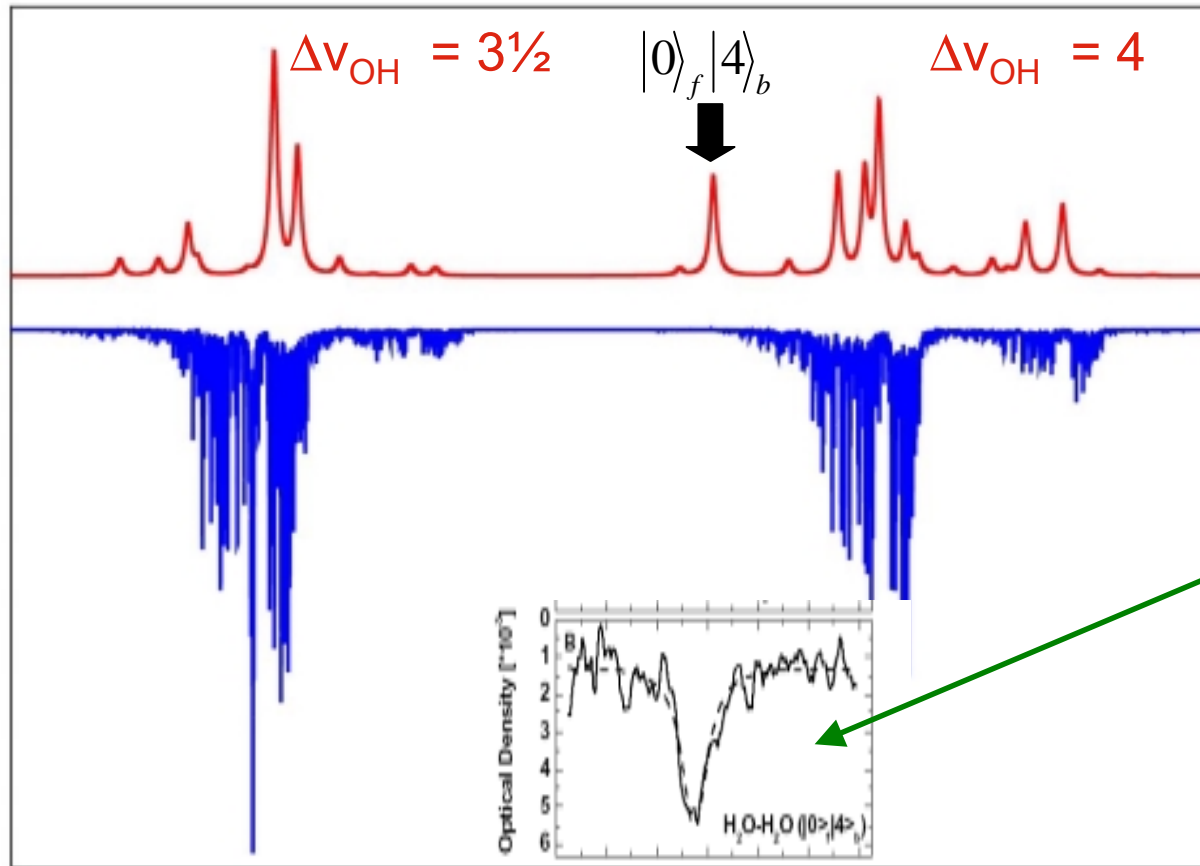
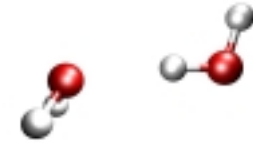
Ptashnik, Smith, Shine, Newnham, QJRMS (2004)



$$|0\rangle_f |1\rangle_b |1\rangle_\beta \quad |10\rangle_- |1\rangle_\beta$$

Schofield, Kjaergaard, PCCP 2003.

Water dimer in the atmosphere



H₂O•H₂O (calc)

H₂O (observed)

Atmos. observation of water dimer.

Abundance ~1/1000 of water monomer.

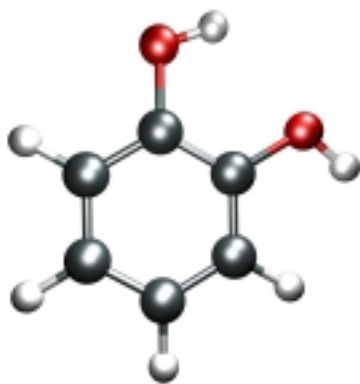
11500

13100
wavenumber (cm⁻¹)

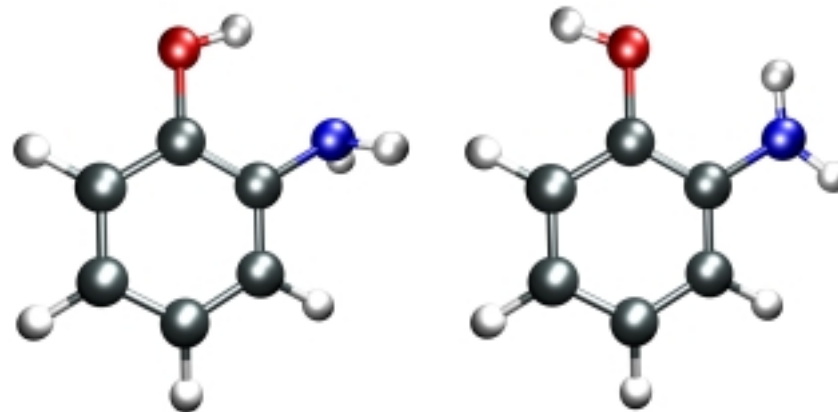
14750

Intra-molecular interaction

Intra-molecular interaction seen in overtone spectra of a few molecules.



Catechol



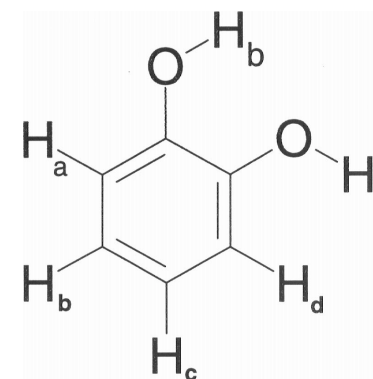
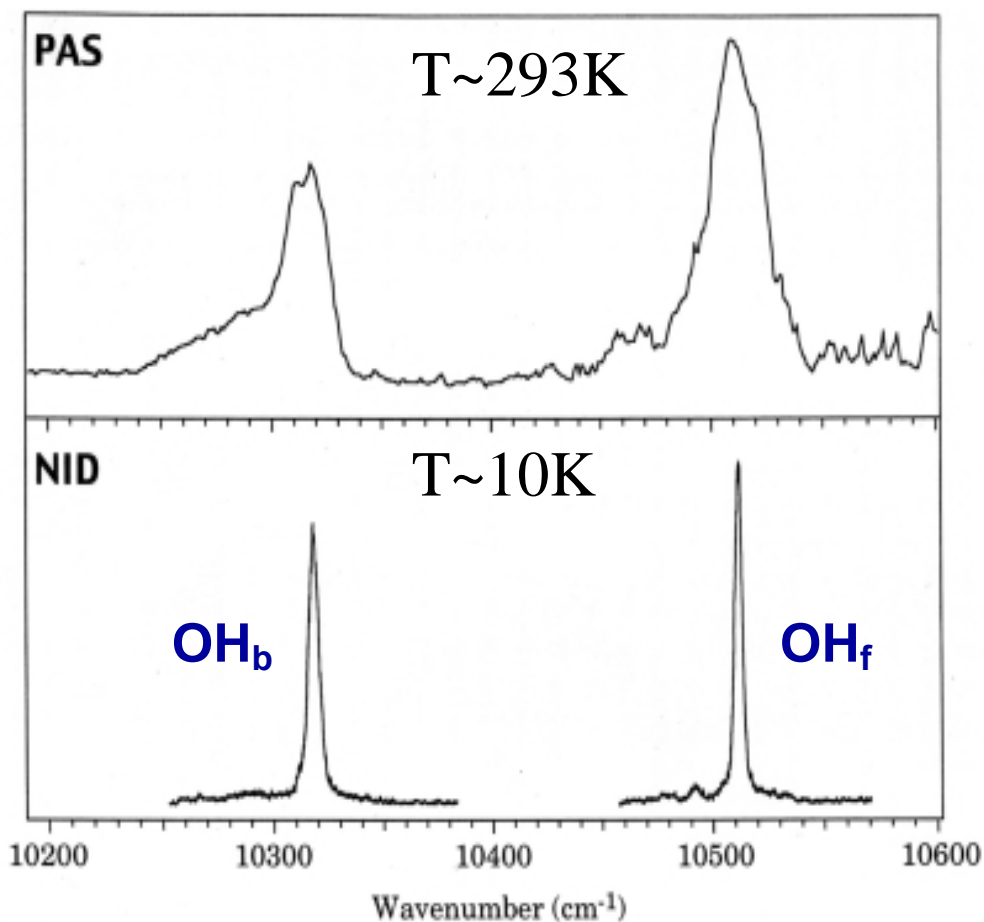
cis- and *trans*-o-aminophenol

Vapor pressure often low and the overtones are inherently weak, so a sensitive technique is necessary.



Catechol

ICL-PAS and NID spectrum of catechol in the $\Delta v_{\text{OH}} = 3$ region



Bond lengths:

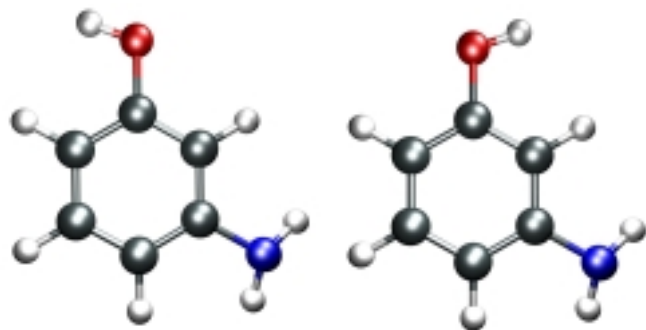
OH_b = 0.964 Å

OH_f = 0.961 Å

B3LYP/6-311++G(2d,2p)

Kjaergaard, Howard, Schofield, Robinson, Ishiuchi, Fujii, JPC A 2002.

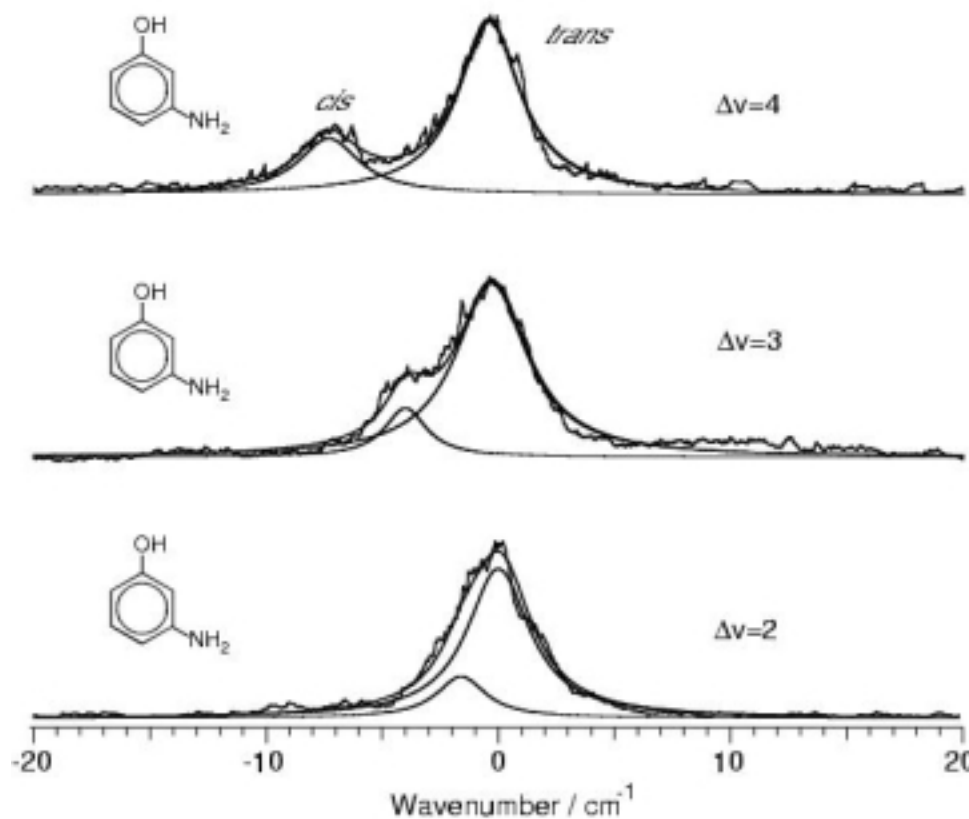
Rotational isomers



cis- and *trans*-*m*-aminophenol.

B3LYP/6-311++G(2d,2p)
Calculated OH bond length
difference is only 0.0002Å

In *o*-aminophenol, the *trans* isomer is observed, as little shift in OH-stretch vs phenol is observed.



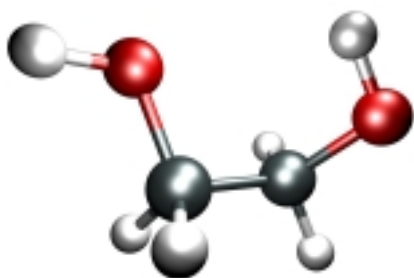
NIR-NIR spectroscopy

Robinson, Kjaergaard, Ishiuchi, Shinozaki, Fujii, JPC A 2004.

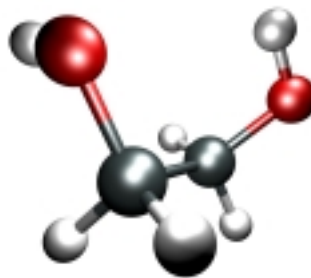
Ethylene Glycol

Different conformers present at room temperature and seen in the overtone spectra.

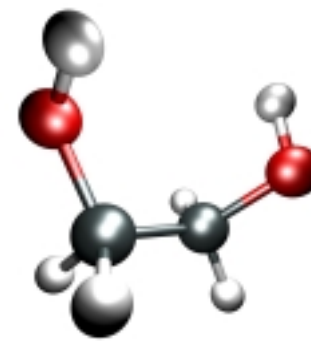
Hydrogen bonding



(~58%)



(~26%)

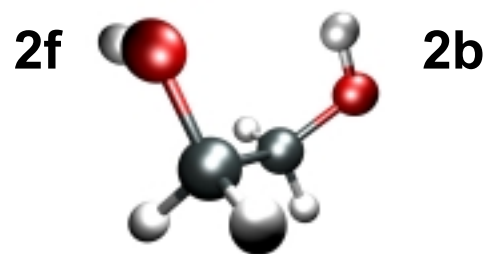
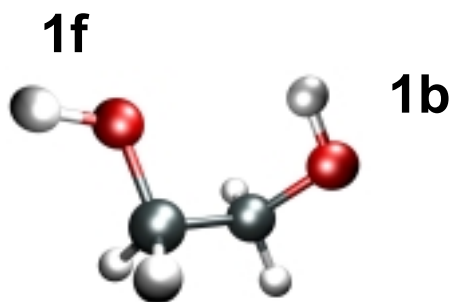


(~10%)

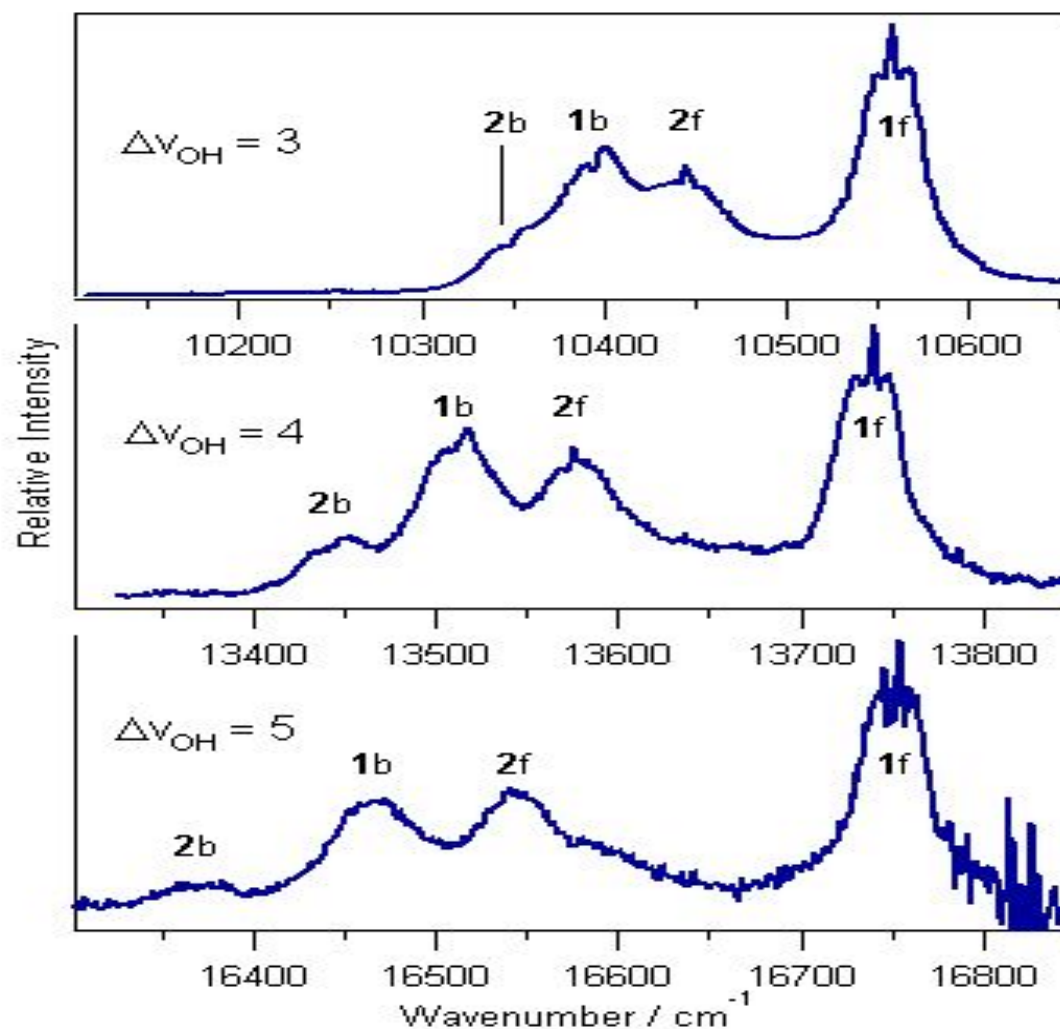
Small molecule, so high level ab initio calculations are possible: CCSD(T)/aug'-cc-pVTZ.

Howard, Jørgensen, Kjaergaard, submitted Aug 05.

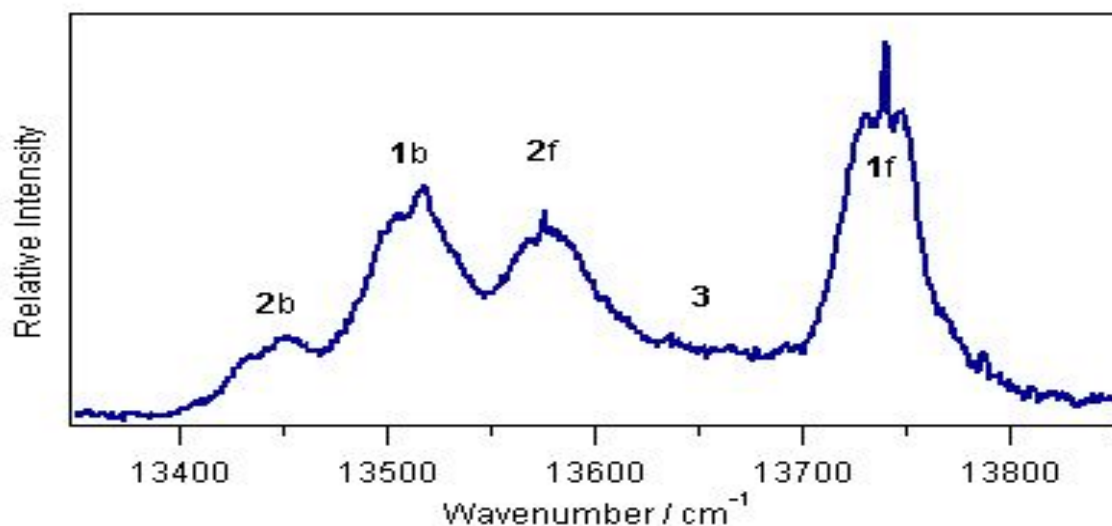
Ethylene Glycol



Higher overtones
better but also more
difficult!



Ethylene Glycol



Conf. 3 predicted in gap $\sim 13666 \text{ cm}^{-1}$.
Supported by temperature study.

The CCSD(T)/aug'-cc-pVTZ calculated frequencies, with no scaling, are close to experimental values.

$$\Delta 1f = +2\text{cm}^{-1}, \Delta 2f = +25\text{cm}^{-1}, \Delta 1b = +56 \text{ cm}^{-1}, \Delta 2b = +43\text{cm}^{-1}$$

**Geometric position of bonds 1b and 2b similar to bond 3.
Both 1b and 2b shifted relative to 3 \Rightarrow Hydrogen bonding!**

Conclusion

- ☺ **Overtone spectroscopy is a very sensitive probe for inter- and intra-molecular interactions.**
- ☺ **We can calculate quite accurate absolute overtone intensities *ab initio* for species that have not been observed.**
 - ☺ **Guide experimental efforts to observe these species.**
 - ☺ **Provide input for atmospheric impact studies.**
- ☺ **The HCAO local mode model give good absolute overtone intensities in neutral molecules. However, care is required with the choice of *ab initio* method for systems with weak inter- or intra-molecular forces**

Acknowledgements

Timothy Robinson
Daniel P. Schofield
Daryl L. Howard
Joseph R. Lane
Anna Garden
Ben Miller

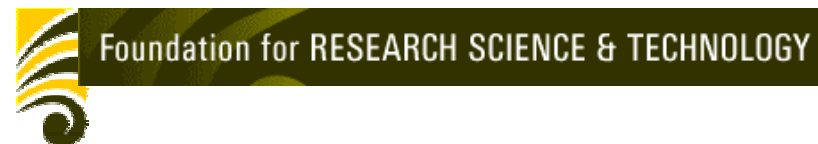
John S. Daniel, NOAA

Bryan R. Henry, Guelph
Veronica Vaida, Boulder
Masaaki Fujii, Tokyo
Terrence Quickenden, Perth
Jean-Pierre Perchard, Paris
Poul Jørgensen, Aarhus



MARSDEN FUND

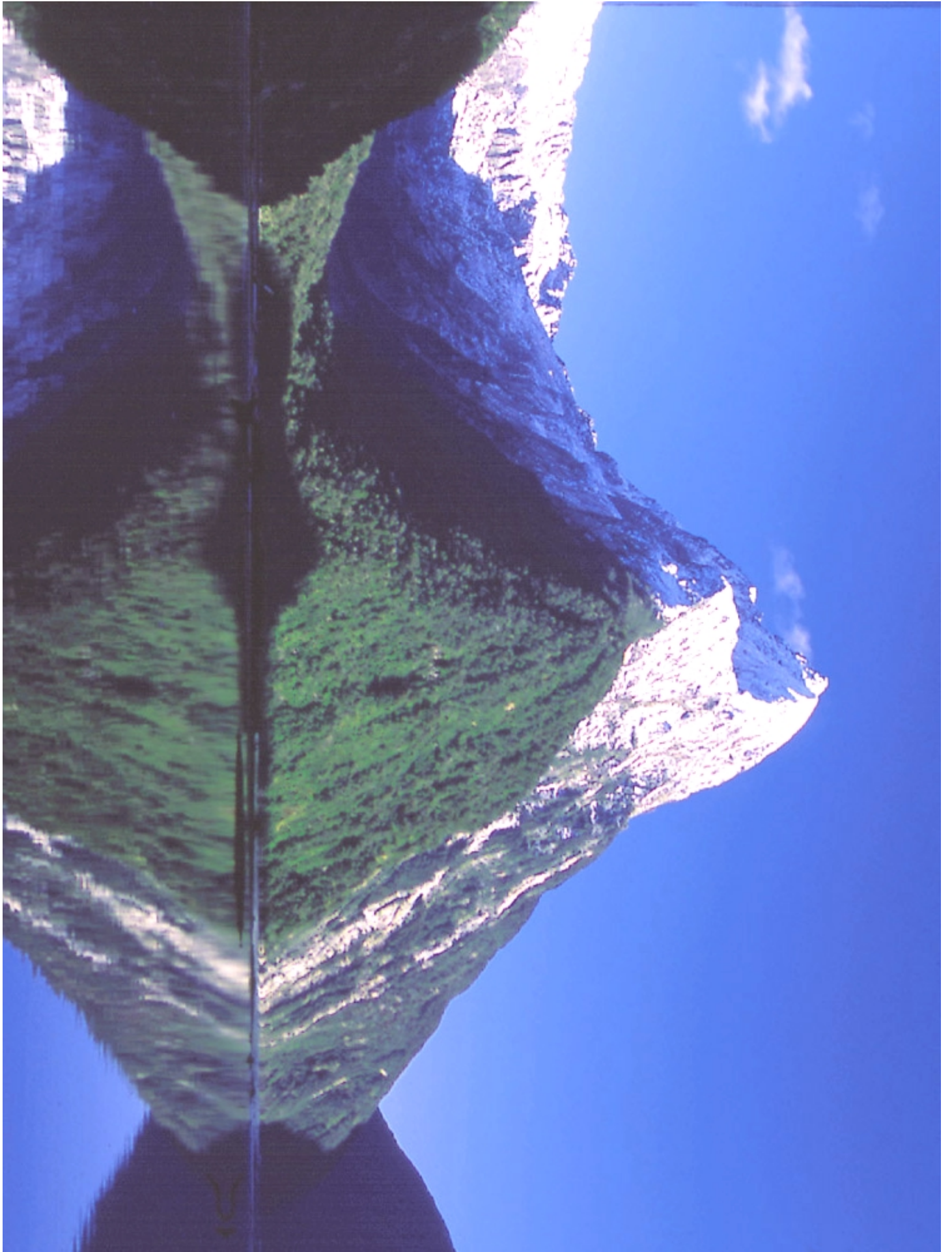
The Royal Society of New Zealand
for the advancement and promotion of science and technology



Durmedin

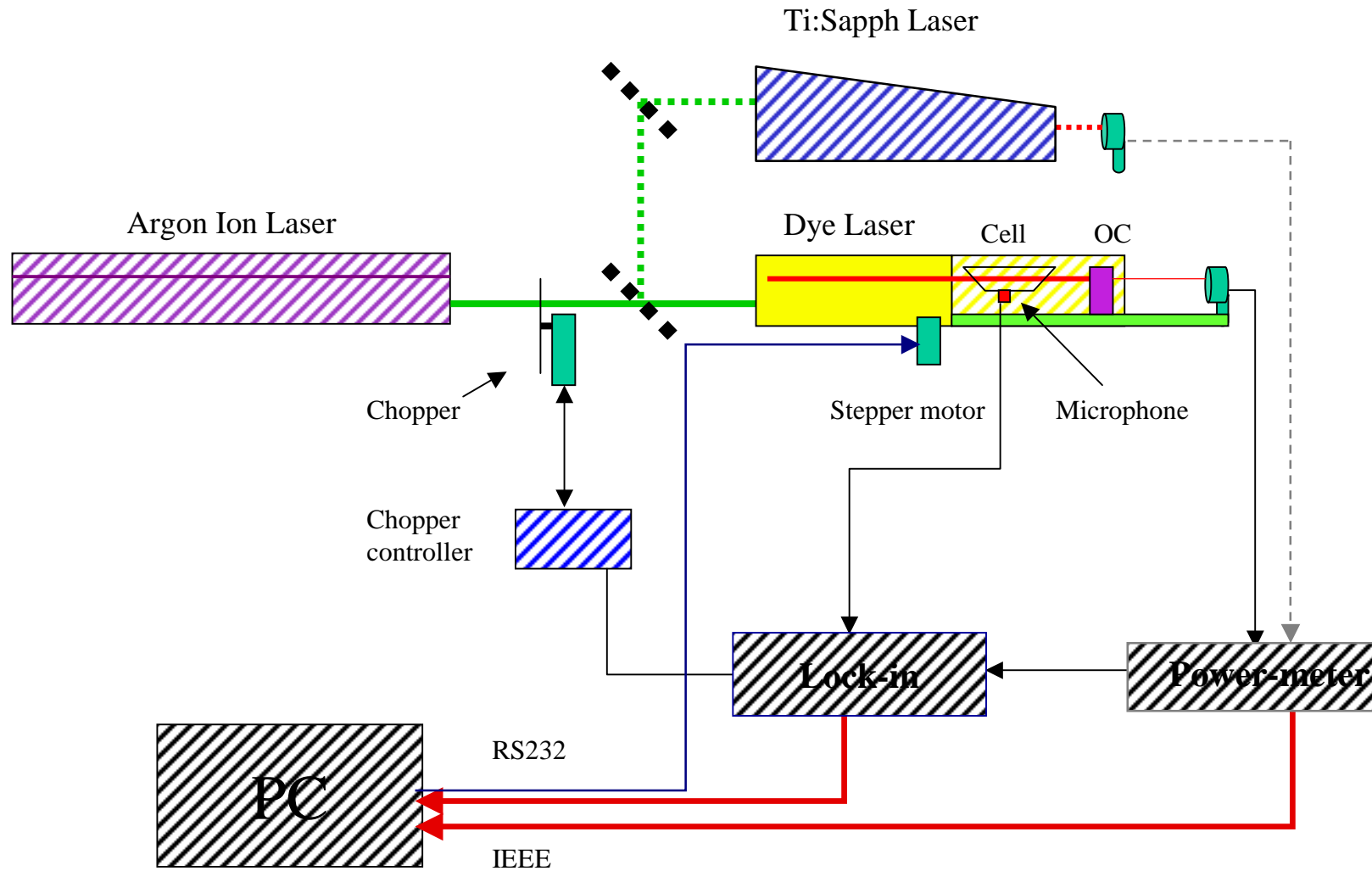






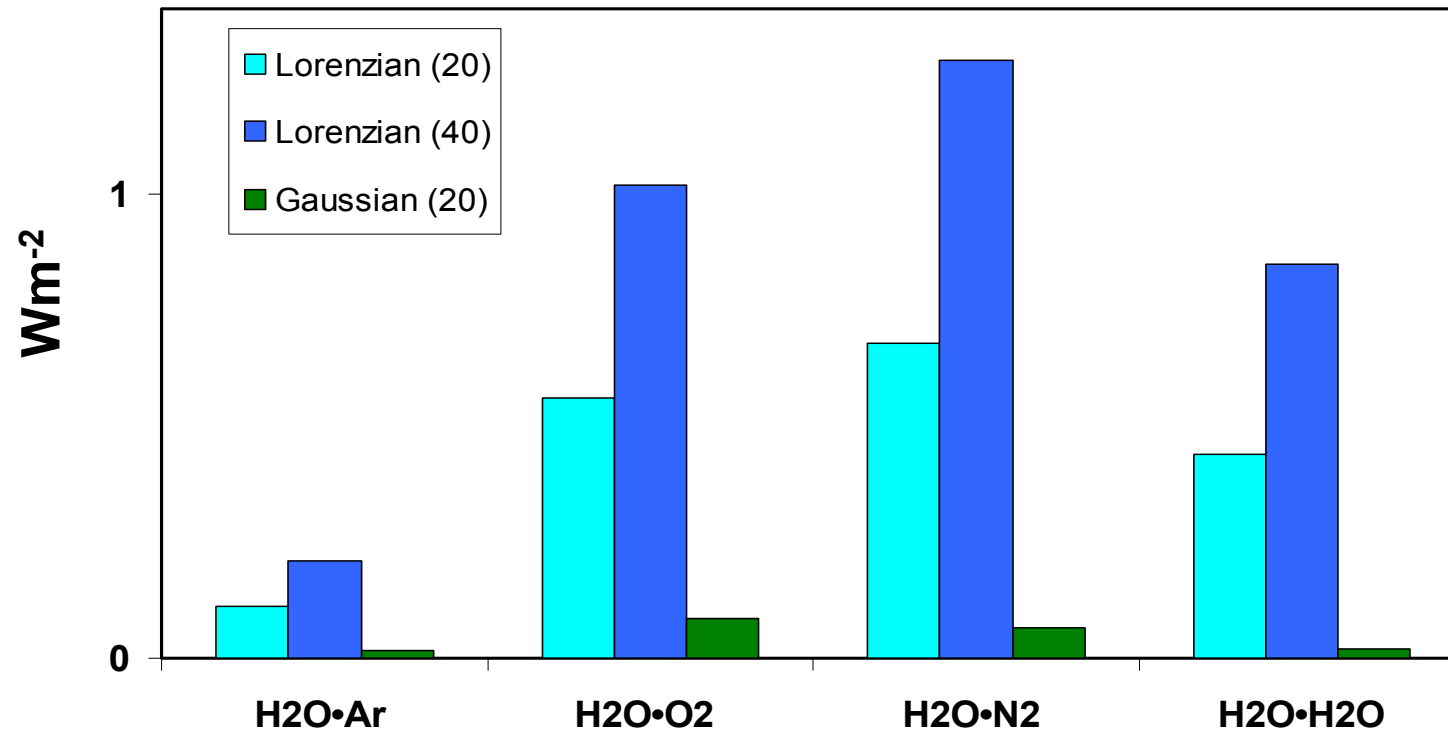


ICL-PAS at Otago



Hydrated complexes

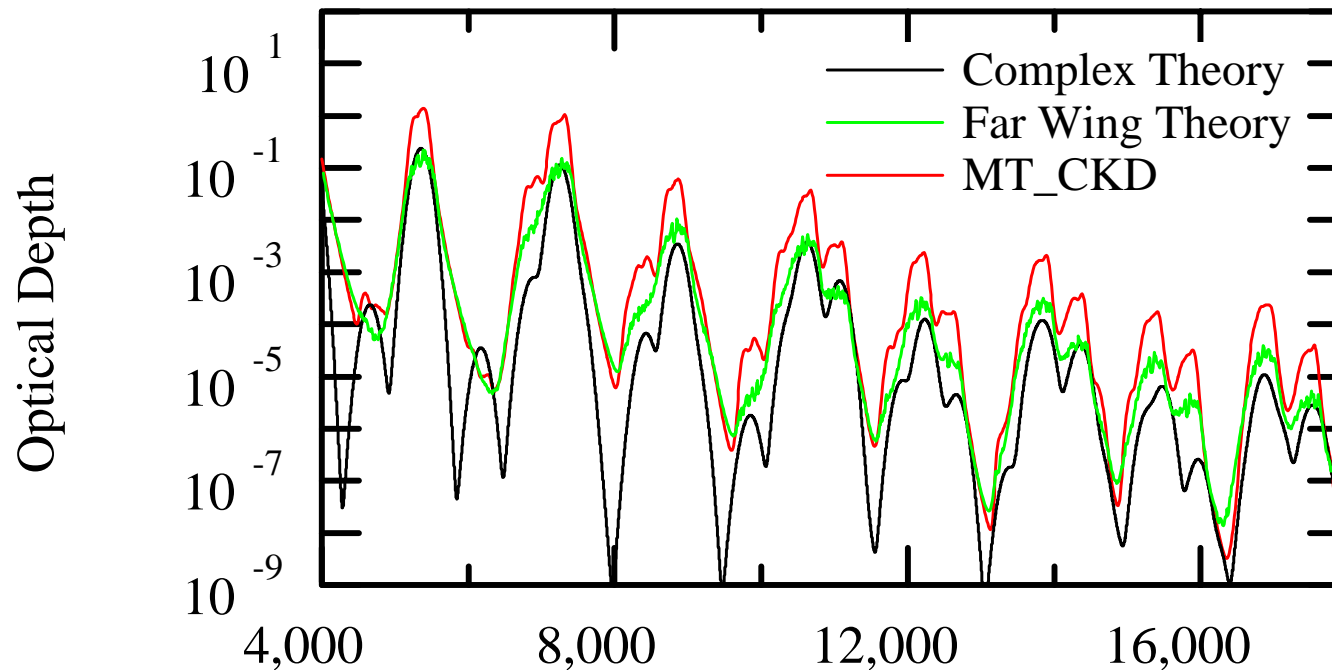
Total absorption of solar radiation



Large dependence on linewidth and shape of the transitions.

Continuum

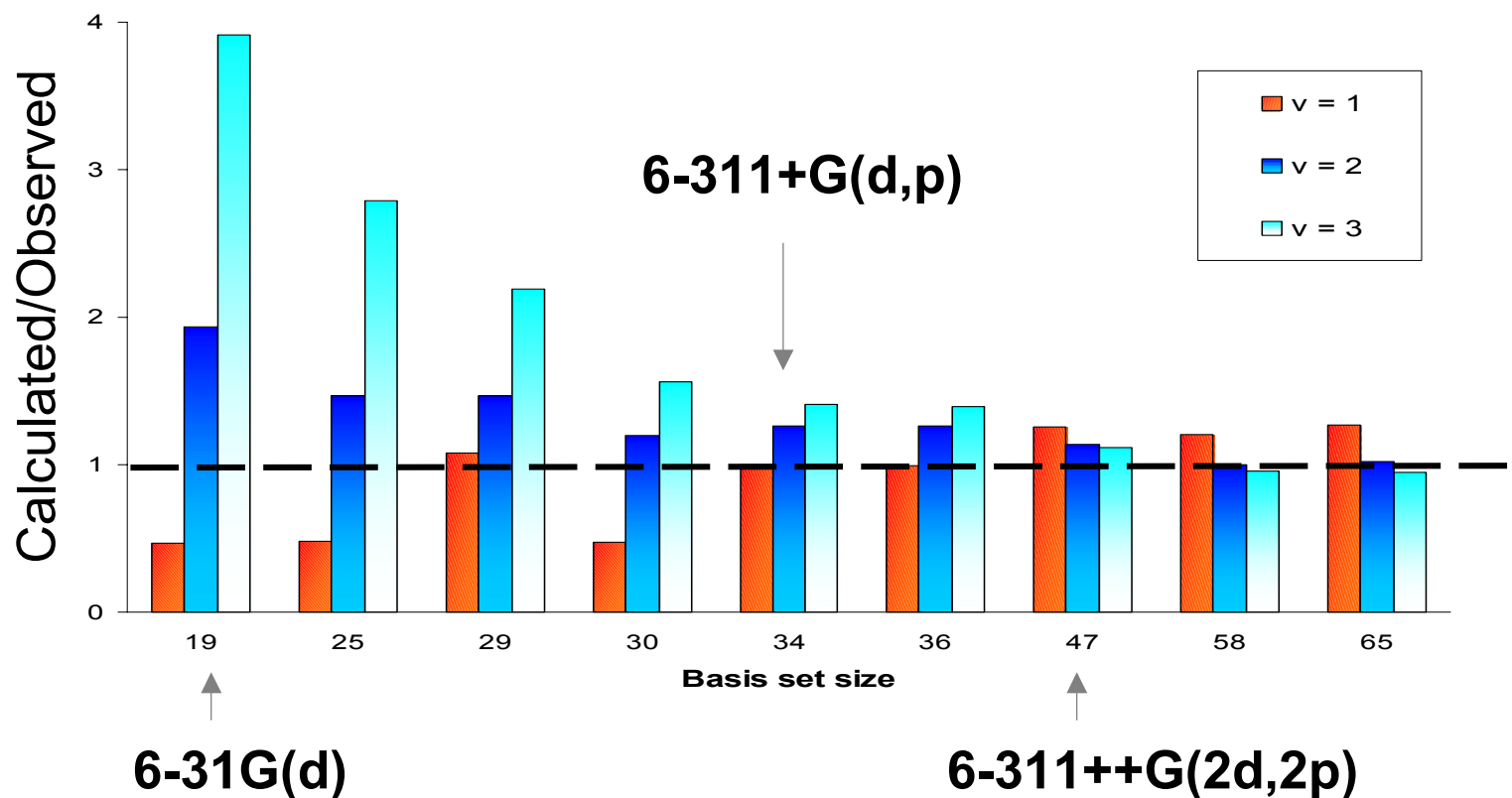
Empirical continuum used to explain discrepancy in models



Seems likely that these hydrated complexes, $\text{H}_2\text{O}\cdot\text{X}$, are the explanation for this continuum

Effect of basis set size

Oscillator strengths of $|v_0\rangle \rightarrow |0\rangle$ states for H₂O

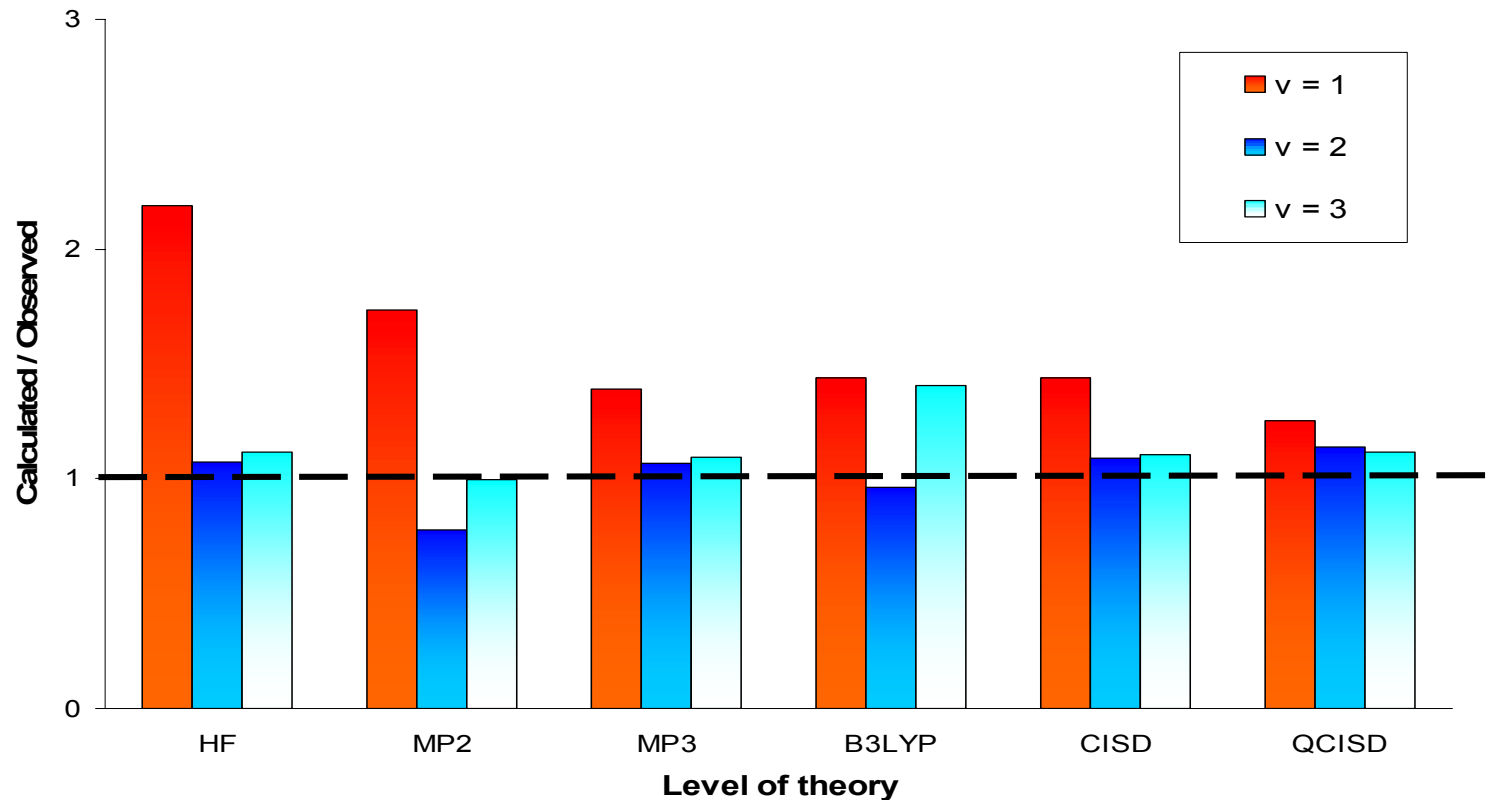


QCISD level of theory & 3D HCAO local mode model

Kjaergaard, Henry, Mol. Phys. 1994.

Effect of level of theory

Oscillator strengths of $|v_0\rangle \rightarrow |0\rangle$ states for H₂O



6-311++G(2d,2p) basis set & 3D HCAO local mode model

HCAO local mode model

Harmonically Coupled Anharmonic Oscillators (HCAO) local mode model. Two equivalent oscillators:

$$\left(\hat{H} - E_{|00\rangle}\right) / hc = (v_1 + v_2) \tilde{\omega} - (v_1^2 + v_1 + v_2^2 + v_2) \tilde{\omega} x - \gamma' (a_1^+ a_2 + a_1 a_2^+)$$

$\tilde{\omega}$, $\tilde{\omega} x$, and γ' from observed peak positions $\Rightarrow E_{vib}$ and ψ_{vib}

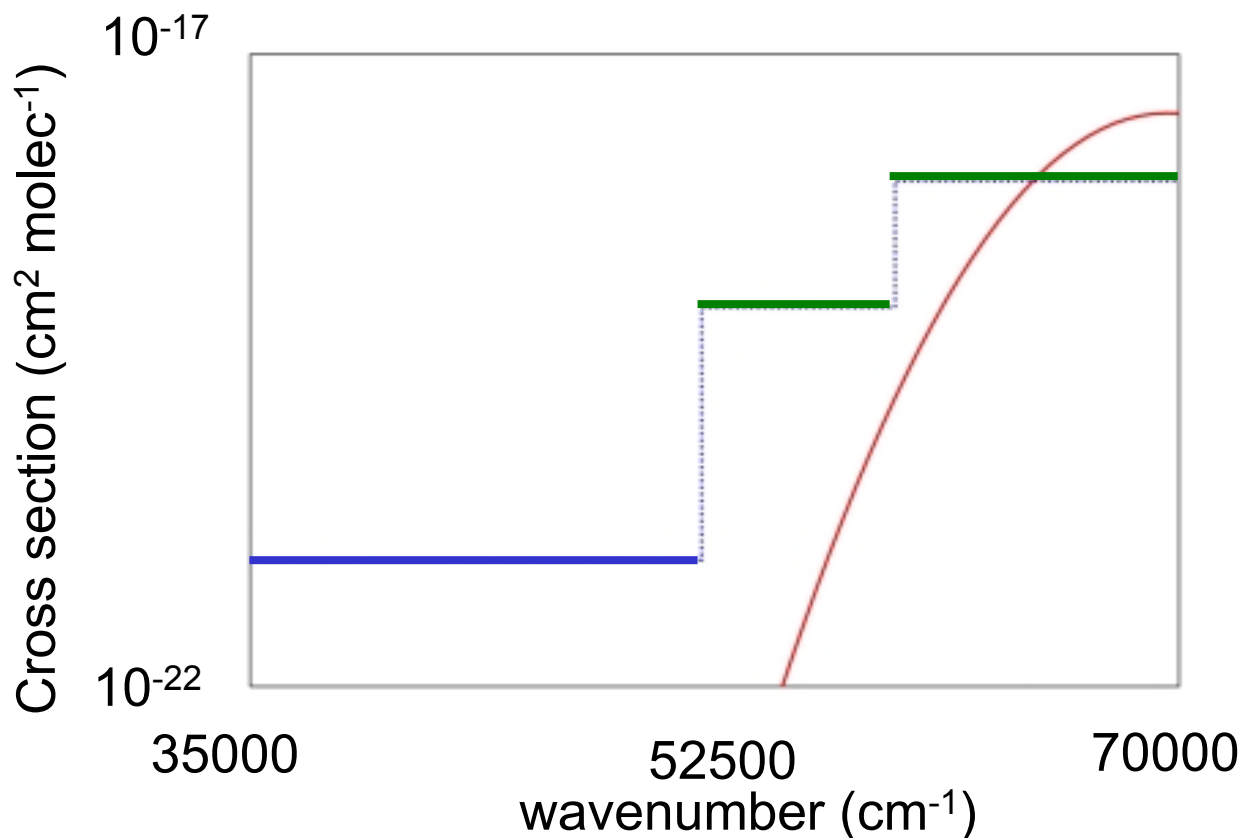
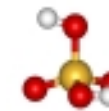
Dipole moment function of two XH bonds

$$\vec{\mu}(q_1, q_2) = \sum_{i,j} \vec{\mu}_{ij} q_1^i q_2^j$$

Coefficients calculated from
2D *ab initio* dipole moment grid

Intensity calculation reduced to 1D integrals like $\langle v | q^n | 0 \rangle$

Electronic spectrum of H₂SO₄



Lowest energy
transition at
 $\lambda = 144 \text{ nm}$
 $f_{\text{calc}} = 0.03$

Solar photons
 $\lambda > 179 \text{ nm}$
 $\sim 55800 \text{ cm}^{-1}$

Burkholder, Mills, McKean, GRL (2000)

Hintze, Kjaergaard, Vaida, Burkholder, JPC A (2003)

Robinson, Schofield,
Kjaergaard, JCP (2003)